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A Single-Molecule Magnet with a "Twist"

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The origin of the ferromagnetic exchange in a recently reported oxo-centered triangular Mn(III) single-molecule magnet (SMM) based on oximato ligands is fascinating since the exchange in all other complexes containing the [Mn^{III}₃O]⁷⁺ core, including the wellknown "basic carboxylates" of general formula [MnIII3O(O2- $CR_{6}L_{3}^{+}$ (R = Me, Et; L = py, MeCN, etc.), is antiferromagnetic.¹ We speculate that a structural distortion in the core of the molecule caused primarily by the "twisting" of the oximato ligands (with regard to the Mn^{III}₃ plane) is responsible. In order to test this hypothesis we have made a family of novel ligands based on salicylaldoxime (saoH₂) in which the oximate carbon atom has been derivatized to possess the "bulky" Me (Me-saoH₂), Et (Et-saoH₂), and Ph (Ph-saoH₂) groups, and synthesized analogues of the known hexanuclear SMM [Mn^{III}₆O₂(sao)₆(O₂CPh)₂(EtOH)₄] (1).² Complex 1, obtained upon the reaction of $Mn(O_2CPh)_2 \cdot 2H_2O$ with saoH₂ in EtOH (Figure S1) contains a nonplanar $[Mn^{III}_{6}(\mu_3-O)_2(\mu-OR)_2]^{12+}$ unit of two off-set, stacked $[Mn^{III}_{3}(\mu_{3}-O)]^{7+}$ triangular subunits linked by two central oximato oxygens, with the remaining four sao²⁻ ligands bridging in a near-planar $\eta^1:\eta^1:\eta^1:\mu$ -fashion along the edges of the $[Mn^{III}_{3}(\mu_{3}-O)]^{7+}$ triangles. The four "central" metals (Mn1, Mn3) are six-coordinate and in distorted octahedral geometries, while the outermost Mn ions (Mn2) have square pyramidal geometries with an additional axial contact of \sim 3.5 Å to a phenolato oxygen. The coordination geometry of the metal ions is completed by a combination of terminal alcohols and μ -carboxylate groups. 1 displays an S = 4 spin ground state as a result of ferromagnetic exchange between the two antiferromagnetically coupled Mn^{III}₃ triangles.²

The idea was to investigate whether the additional steric bulk of the derivatized oximates would enforce structural distortions similar to those seen in the ferromagnetic Mn₃ triangle¹ and the recently obtained ferromagnetic [Mn^{III}₄(Me-sao)₄(Me-saoH)₄] "cube",³ and thus switch the magnetic behavior from antiferromagnetic to ferromagnetic. Complex [Mn^{III}₆O₂(Et-sao)₆(O₂CPh)₂(EtOH)₄(H₂O)₂]. 2EtOH (2·2EtOH) which contains the ethyl-derivatized saoH₂ (Et $saoH_2$) does indeed display ferromagnetic exchange (S = 12 ground state). Here we discuss its structure and magnetic properties. Complex 2 also crystallizes in the triclinic space group P1 and its structure (Figure 1) is analogous to that of 1.4 However, the increased steric bulk of the Et-sao²⁻ ligands causes a shortening of the phenolato oxygen (O71)-square pyramidal Mn (Mn3) distance (~2.5 Å) and severe twisting of the Mn-N-O-Mn moieties within each Mn₃ subunit (Figure 1). This is evidenced by the average Mn–N–O–Mn torsion angle, which in 1 is $\alpha_v = 17.5^{\circ}$ compared to $\alpha_v = 36.5^\circ$ for 2. This also results in a change in the



Figure 1. The molecular structure of complex 2, highlighting its core (top); the "twisted" Mn-N-O-Mn arrangement of the Etsao2- ligands in 2 (bottom).

coordination of the carboxylates: from μ -bridging to terminal, with the "vacant" site now occupied by an additional solvent (alcohol) molecule.

Variable-temperature dc magnetic susceptibility data were collected on 2 in the temperature range 5-300 K in an applied field of 0.1 T (Figure S2). The room temperature $\chi_M T$ value of 19.24 cm³ K mol⁻¹ is slightly above that expected for six noninteracting Mn^{III} ions (18 cm³ K mol⁻¹). Upon cooling the value of $\chi_M T$ increases gradually to 23.72 cm³ K mol⁻¹ at \sim 70 K, below which it increases rapidly to a maximum value of 69.57 cm³ K mol⁻¹ at 6.5 K, before decreasing slightly to 69.36 cm³ K mol⁻¹ at 5 K. This behavior is indicative of ferromagnetic exchange between the metal centers with the low-temperature maximum suggesting an S = 12 ground state. In order to confirm the ground state of the molecule, magnetization data were collected in the ranges 0.5-7 T and 1.8-7 K, and these are plotted as reduced magnetization $(M/N\mu_B)$ versus H/T in Figure 2. The data were fit by a matrix diagonalization method to a model that assumes only the ground state is populated, includes axial zero-field splitting $(D\hat{S}_z^2)$ and the Zeeman interaction, and carries out a full powder average. The corresponding Hamiltonian is given by eq 1,

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Figure 2. Plot of reduced magnetization $(M/N\mu_B)$ versus H/T for **2** in the field and temperature ranges 0.5–7 T and 1.8–7 K. The solid lines correspond to the fit of the data.



Figure 3. Out-of-phase (χ_M'') ac susceptibility measurements in the 1.8–8 K and 50–1000 Hz ranges (insert); Arrhenius plot using ac (green) and dc (red) data. The dashed line is the fit of the thermally activated region to eq 2.

$$\hat{H} = D\hat{S}_z^2 + g\mu_{\rm B}\mu_0\hat{S}\cdot H \tag{1}$$

where *D* is the axial anisotropy, $\mu_{\rm B}$ is the Bohr magneton, μ_0 is the vacuum permeability, \hat{S}_z is the easy-axis spin operator, and *H* is the applied field. The best fit gave S = 12, g = 1.99 and D = -0.43 cm⁻¹.

Ac susceptibility measurements were performed in the 1.8–16 K range under a 3.5 G ac field oscillating at 50–1000 Hz. The value of the in-phase $(\chi_M'T)$ signal (Figure S3) increases with decreasing temperature to ~6 K where it then plateaus at ~4 K before displaying a frequency-dependent decrease below this temperature. This is indicative of the presence of excited states of smaller *S* values relatively close to the ground state. Extrapolation of the in-phase $(\chi_M'T)$ signal to 0 K from the plateau region gives a value of ~74 cm³ K mol⁻¹ indicative of an *S* = 12 ground state, in agreement with the dc data (Figure S3). Fully visible, frequency-dependent out-of-phase (χ_M'') ac susceptibility signals are seen below ~6 K with the peak at 1000 Hz occurring at ~4.5 K (Figure 3).

Single-crystal hysteresis loop and relaxation measurements were performed on **2** using a micro-SQUID setup.⁵ Hysteresis loops were observed whose coercivity was strongly temperature (Figure 4) and sweep-rate (Figure S4) dependent, with a blocking temperature of \sim 3.5 K. Relaxation data determined from a combination of single-crystal dc and powder ac measurements were fitted to the Arrhenius relationship (eq 2, Figure 3),

$$\tau = \tau_0 \exp(U_{\text{eff}}/kT) \tag{2}$$

where $U_{\rm eff}$ is the effective relaxation barrier, τ is the relaxation time, and k is the Boltzmann constant, giving $U_{\rm eff} = 53.1$ K and $\tau_0 = 8 \times 10^{-10}$ s. The theoretical upper limit of $U = S^2 |D| = 88.5$ K is strongly reduced by the presence of low-lying excited states as



Figure 4. Magnetization versus field hysteresis loops for a single crystal of **2** at the indicated temperatures in a field sweep rate of 0.14 T s⁻¹. *M* is normalized to its saturation value.

evidenced by the fine structure in the hysteresis loops which are indicative of tunneling between excited-state multiplets. Detailed analyses of these will appear in a full paper.

In conclusion, the deliberate structural distortion of a Mn_6 compound via the use of a bulky salicylaldoxime derivative switches the intra-triangular magnetic exchange from antiferromagnetic to ferromagnetic resulting in an S = 12 ground state and an anisotropy barrier (U_{eff}) approaching that of the Mn_{12} family.^{7,8} The combined results thus suggest that the intelligent use of ligands that can cause targeted distortions to the core structures of Mn_6 (and other) clusters may prove to be a viable means to raise the blocking temperatures of SMMs to record values.

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Supporting Information Available: Crystallographic details in CIF format; synthetic procedures: and magnetism data. This material is available free of charge via the Internet at http://pubs.acs.org.

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